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Estimated strength of Atlantic overturning circulation during the last

deglaciation

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The reconstruction of the state and strength of the Atlantic meridional overturning circu-1 lation (AMOC) in the past is an important topic in paleoceanography. Many proxies have 2 been proposed that give information on the state of the past circulation¹. Of these, only the 3 protactinium-to-thorium (²³¹Pa/²³⁰Th) activity ratio of ocean sediments and top-to-bottom ra-4 diocarbon age differences provide qualitative estimates of changes in the AMOC strength^{2,3,4,5}. But the robustness of these indicators is still under debate 6,7,8 . Here we reconstruct changes in the AMOC strength over the period of the last deglaciation quantitatively by using a novel 7 method, which derives changes in the ocean circulation from a combination of time series re-8 constructions of sea surface temperature (SST) and air temperature. A global climate model 9 provides the relationship between SST, air temperature, and AMOC strength. Our recon-10 struction of the changes in the AMOC strength during the last deglaciation is consistent with 11 qualitative reconstructions from sedimentary 231 Pa $/^{230}$ Th. It suggests an AMOC reduction 12 of about 14 Sv (Sverdrup; 1 Sv = 10^6 m³ s⁻¹) during Heinrich stadial 1 compared to the Last 13 Glacial Maximum (LGM), and a slightly lower reduction during the Younger Dryas cold event. 14 The difference between the AMOC strength during the LGM and the Holocene lies within 15 the uncertainty of the reconstruction. 16

The AMOC is an important feature of the Earth's climate system. It is characterized by northward flowing 17 water masses at the surface ocean, deep water formation in the North Atlantic and southward flowing water 18 masses in the intermediate to deep ocean. This overturning cell is complemented by an underlying, generally 19 weaker, reversed overturning cell that originates in the Southern Ocean. In the following, we use AMOC 20 to refer only to the upper overturning cell. Because of the large heat capacity of water, the northward 21 flowing near-surface waters contribute effectively to the heat transport from the tropics to the mid- and 22 high latitudes and therefore affect the regional climate around the North Atlantic^{9,10}. Hence, changes of 23 the AMOC strength can substantially influence climate in the north, especially in Europe. The depth of 24 the southward flowing waters as well as the strength of the circulation have changed in the past. During 25 the last glacial period, abrupt warming events of several degrees within a few decades followed by slower 26 cooling were detected in Greenland ice cores¹¹. Because effects of the so-called Dansgaard-Oeschger events 27

were found far beyond the Greenland ice sheet ^{12,13,14} with patterns reminiscent of varying inter-hemispheric heat distribution ^{15,16}, they are believed to be the result of rapid changes of the AMOC strength ¹⁷. Other examples of a reduction or even a complete shutdown of the AMOC are Heinrich stadial 1 from approximately 16.8 to 14.6 ka before present (B.P., before year 1950 AD) ¹⁸, the Younger Dryas stadial from approximately 12.8 to 11.6 ka B.P. ¹⁹ and the 8.2 ka event ²⁰.

In this study, we present a numerical reconstruction scheme to infer the past evolution of the AMOC 33 strength on centennial to glacial-interglacial timescales. The scheme is based on earlier work²¹ that proposes 34 that ocean temperature anomalies at any location (i), $T_{\rm oc}^{\prime i}$, can be approximated by a linear combination 35 of anomalies of global mean air temperature $(T'_{\rm atm})$, AMOC $(\Psi'_{\rm AMOC})$, and Southern Ocean meridional 36 overturning circulation (SOMOC) strength. The strength of the circulation is defined by the maximum of 37 the meridional overturning stream function. In this study we focus on the surface Atlantic, where the effect 38 of the SOMOC can be neglected 21 . Hence, we can linearly approximate the time evolution of the surface 30 temperature pattern $T_{\rm oc}^{\prime i}(t)$ as 40

$$T_{\rm oc}^{\prime i}(t) = a_{\rm Tatm}^{i} \cdot T_{\rm atm}^{\prime}(t) + a_{\rm AMOC}^{i} \cdot \Psi_{\rm AMOC}^{\prime}(t).$$
(1)

⁴¹ The relationships between Atlantic SST, air temperature and AMOC are the result of ocean dynamics, and ⁴² we determine coefficients (a^i_{Tatm} and a^i_{AMOC}) describing the postulated linear relationships using a global ⁴³ climate model, subject to a combination of forcings (see Methods section). Global mean air temperature ⁴⁴ time series can be determined from air temperature reconstructions from Antarctic ice cores ²² and SST time ⁴⁵ series reconstructions are available from several marine sediment cores. Hence, Eq. (1) can be solved for ⁴⁶ Ψ'_{AMOC} . When several SST records are incorporated into the reconstruction, a least squares solution can be ⁴⁷ calculated:

$$\Psi'_{\rm AMOC}(t) = \sum_{i=1}^{n} \frac{w_i}{a^i_{\rm AMOC}} [T'^i_{\rm oc}(t) - a^i_{\rm Tatm} T'_{\rm atm}(t)],$$
(2)

where w_i is the weight of location *i* and is determined by the least squares solver.

The simplicity of our assumption requires thorough testing of the reconstruction method. This is accomplished by hindcasting the AMOC strength of two model simulations from two independent climate models.

Also, the coefficients of the reconstruction scheme are calculated by two independent models (Fig. 1): By 51 an ensemble of simulations performed with the Bern3D Earth System Model of intermediate complexity 52 (shorthand B3D; Supplementary information; Fig. S2; Table S3), and by a transient deglaciation simulation 53 performed using the LOVECLIM Earth System Model of intermediate complexity²³ (shorthand LVC; Sup-54 plementary information; Fig. S4; Table S4). In both cases a broad range of climate states is generated in 55 response to various freshwater and CO₂ forcings that alter the AMOC, SST and global mean air tempera-56 ture, respectively. Using the coefficients derived from the B3D and the LVC sensitivity experiments, AMOC 57 hindcasts are calculated for the LVC simulation, and for a freshwater forcing simulation conducted with 58 the NCAR CSM1.4 coupled atmosphere-ocean general circulation model²⁴ (shorthand CSM; Supplementary 59 information; Fig. S5). 60

For the AMOC hindcasts, we substitute the real world by one of the model simulations and "reconstruct" the AMOC evolution based on our method using simulated SST time series from several Atlantic locations as well as global mean air temperature time series. The reconstructed AMOC evolution can then be compared to the temporal evolution of the true model AMOC. For both AMOC hindcasts, 14 model SST time series of the Atlantic Ocean are used (Fig. S1). The locations of the time series correspond to locations where real SST reconstructions of the last deglaciation are available with an average resolution of at least 500 years (Table S1). These SST records will be used later for the reconstruction of the deglacial AMOC variability.

As a first test, an AMOC hindcast of the LVC simulation is calculated, where the reconstruction scheme 68 coefficients are also determined from the LVC simulation (Fig. 2a). The hindcast correlates well with the 69 modeled AMOC (R = 0.97) and the relative standard deviation of $s_{\rm rel} = 1.03$ is close to the optimum value 70 of 1 ($s_{\rm rel}$ is defined as the standard deviation of the hindcasted AMOC divided by the standard deviation 71 of the modeled AMOC). The LVC hindcast is also successful when coefficients determined from the B3D 72 simulations are used $(R = 0.97, s_{rel} = 0.85;$ Fig. 2a). The largest discrepancy between hindcast and model 73 AMOC appears during the Bølling between 15 and 14 ka B.P., where the AMOC peak of the model is not 74 captured by the hindcast. The peak would be somewhat better captured if SSTs from the central North 75 Atlantic (at approximately 50°N, 35°W) were incorporated into the hindcast (Fig. S6). Unfortunately, there 76 are currently no paleo-SST reconstructions available from this region and hence we have not included this 77

⁷⁸ region in our analysis.

The weights that the least squares solver assigns to the individual SST locations determines the relative contributions of the individual SST time series to the AMOC hindcast. The smaller the weight, the less sensitive the temperature at the particular location is to changes of AMOC strength. The six northeastern Atlantic SST time series together have the largest influence on the hindcasts and account for a cumulative weight of ca. 90 % (see Table S5 for a list of the weights).

⁸⁴ Next, the AMOC of the CSM model is hindcasted (Fig. 2b). The atmospheric component of this model is ⁸⁵ considerably more complex than LVC. Again, both hindcasts calculated using LVC and B3D coefficients are ⁸⁶ very similar and compare well with the CSM model AMOC (for LVC coefficients: R = 0.98, $s_{rel} = 1.12$; for ⁸⁷ B3D coefficients: R = 0.97, $s_{rel} = 0.96$).

The hindcasts of the LVC deglaciation simulation and of the CSM freshwater forcing simulation emphasize 88 the robustness of the AMOC reconstruction method. In the following, we apply the method to Atlantic SST 89 reconstructions from marine sediment cores to estimate the AMOC anomalies of the last deglaciation. One 90 of the 14 SST records only goes back to 19 ka B.P. We do not use this record for the AMOC reconstruction 91 in order to extend the reconstruction back to the LGM. Because of the small weight of this location, this 92 has little impact on the result. The SST data are splined using a cutoff period of $\tau_{\rm cut} = 1000$ years. Since 93 there is no direct proxy for global mean air temperature, it is determined from the Antarctic air temperature 94 reconstruction from the EPICA Dome C ice $core^{25}$ by lowpass filtering the data using a cutoff period of 10 95 ka to remove local-to-hemispheric scale variability (such as the Antarctic Cold Reversal) and by applying 96 a polar amplification factor of 2 which is determined from comprehensive global atmosphere-ocean general 97 circulation models from the Paleoclimate Modelling Intercomparison Projects PMIP-1 and PMIP-2²². 98

⁹⁹ Two AMOC reconstructions are presented: one using coefficients calculated from the B3D simulation, and ¹⁰⁰ one using coefficients calculated from the LVC simulation (Fig. 3). The 1 σ -uncertainty band of the AMOC ¹⁰¹ reconstruction follows from Gaussian uncertainty propagation of the individual uncertainties (Supplementary ¹⁰² information). These are the uncertainties of the SST reconstructions (Table S1) and the air temperature ¹⁰³ reconstruction ($\pm 1.5^{\circ}$ C), the uncertainties of the age scales of the individual temperature reconstructions $(\pm 500 \text{ yr} \text{ for the last deglaciation})$, and the uncertainties of the coefficients of the reconstruction scheme. The uncertainty band is calculated separately for the reconstruction using the B3D coefficients and for the reconstruction using the LVC coefficients.

The two reconstructions are very similar (Fig. 3). The AMOC during Heinrich stadial 1 (H1) appears to be 107 approximately 14 Sv weaker than during the LGM with a range of 5 to 23 Sv decrease. With a present-day 108 AMOC strength of around 19 Sy ± 5 Sy²⁶, this suggests a substantial slowdown of the AMOC during H1. 109 The reconstruction also shows an AMOC slowdown during the Younger Dryas (YD) by 12 Sv with a range 110 of 5 to 22 Sv decrease in circulation compared to the Bølling/Allerød. The reconstruction is consistent with 111 the AMOC reconstruction by McManus et al². However, our reconstruction suggests a stronger reduction of 112 the AMOC during the YD. Furthermore, the LGM-to-Holocene AMOC difference is less prominent in our 113 reconstruction than in the 231 Pa/ 230 Th data. 114

¹¹⁵ We can use our method to inform about locations for SST records which would be most effective for an ¹¹⁶ AMOC reconstruction. Eq. (2) can be rearranged to

$$\Psi'_{\text{AMOC}}(t) = \sum_{i=1}^{n} \frac{w_i}{a_{\text{AMOC}}^i} T'_{\text{oc}}^i(t) - \beta T'_{\text{atm}}(t), \qquad (3)$$

where $\beta = \sum_{i=1}^{n} (w_i a_{\text{Tatm}}^i / a_{\text{AMOC}}^i)$, and w_i / a_{AMOC}^i is the spatial fingerprint of the AMOC. It determines the influence of every location (*i*) on the AMOC hindcast. To assess these fingerprints, we calculated an AMOC hindcast of the CSM simulation based on SST time series of every LVC surface Atlantic model cell and by using reconstruction scheme coefficients determined from the LVC model (Fig. 4). Not surprisingly, the result is a bipolar seesaw pattern with largest fingerprints in the North East Atlantic.

Our AMOC reconstruction method can also be applied on longer timescales to reconstruct the AMOC over multiple glacial-interglacial cycles. However, a substantially larger age-scale uncertainty must be applied when radiocarbon dating is not available. Furthermore, changes of other large scale circulations can potentially be reconstructed such as the Southern Ocean meridional overturning circulation (SOMOC). In this case, deep ocean temperature time series reconstructions would be suitable candidates for the reconstruction. If deep ocean temperature reconstructions are included, then time lags between changes in air temperature, ocean circulation, and ocean temperatures need to be allowed for in Eq. (1), i.e. $T'_{\text{atm}}(t - \tau^i_{\text{Tatm}})$ and $\Psi'_{\text{SOMOC}}(t - \tau^i_{\text{SOMOC}})$. Further, a possible occurrence of a North Pacific meridional overturning circulation (NPMOC) cell during Heinrich stadials could be reconstructed ²⁷. However, both SOMOC and NPMOC reconstructions would require further ocean temperature time series reconstructions at locations that are sensitive to these circulation changes.

133 Methods

134 Methodology

The reconstruction of the AMOC strength is based on the assumption that on centennial to glacial-interglacial timescales, SST anomalies $(T_{\rm oc}^{\prime i})$ at any location in the Atlantic *i* can be approximated by a linear combination of global mean air temperature anomalies $(T_{\rm atm}^{\prime})$ and AMOC strength anomalies $(\Psi_{\rm AMOC}^{\prime})$:

$$T_{\rm oc}^{\prime i}(t) = a_{\rm Tatm}^{i} \cdot T_{\rm atm}^{\prime}(t) + a_{\rm AMOC}^{i} \cdot \Psi_{\rm AMOC}^{\prime}(t).$$

We define the AMOC strength as the maximum of the Atlantic meridional overturning stream function. Anomalies are defined as $T_{\rm oc}^{\prime i}(t) = T_{\rm oc}^{i}(t) - \overline{T}_{\rm oc}^{i}$ with $T_{\rm oc}^{i}(t)$ the SST at time t and $\overline{T}_{\rm oc}^{i}$ the temperature averaged over a period of time, and analogously for $T_{\rm atm}^{\prime}$ and $\Psi_{\rm AMOC}^{\prime}$.

This approximation is supported by a recent study²⁸ in which a principal component analysis of globally distributed paleo-records of temperature revealed that 61 % of the climate variability of the last deglaciation can be explained by a globally near-uniform pattern that is strongly correlated to the atmospheric CO₂ change (representing global mean $T'_{\rm atm}$), and 11 % can be explained by the bipolar seesaw pattern between the hemispheres¹⁶ which resembles the millennial-scale variability of the AMOC.

In this reconstruction scheme we also assume that coupled climate models are able to provide a reasonable physical relationship between changes of AMOC, air temperatures and ocean temperatures. Accepting these assumptions, then the coefficients of the reconstruction scheme, a_{Tatm}^i and a_{AMOC}^i , can be diagnosed from a ¹⁴⁹ global coupled ocean-atmosphere climate model. Then with $T'_{\rm atm}$ time series reconstructions from ice cores ¹⁵⁰ and $T'_{\rm oc}$ time series reconstructions from marine sediment cores, Eq. (1) can be solved for $\Psi'_{\rm AMOC}$.

A model simulation is performed where a broad range of climate states is produced by subjecting the model to various types of transient perturbations, e.g., by adding and removing freshwater from the North Atlantic to alter the AMOC strength, and e.g., by adding and removing CO₂ from the atmosphere to change air temperature. With this simulation we produce artificial climate reconstructions, i.e. purely model-based time series of $T_{oc}^{\prime i}$, T_{atm}^{\prime} and Ψ_{AMOC}^{\prime} . At the location of every SST reconstruction (*i*), a_{Tatm}^{i} and a_{AMOC}^{i} are determined by multiple linear regression of the simulated SST evolution to the reconstruction scheme using the simulated global mean air temperature as well as the simulated AMOC strength:

$$(T_{\rm oc}^{\prime i})^{\rm fit}(t) = a_{\rm Tatm}^{i} \cdot (T_{\rm atm}^{\prime})^{\rm model}(t) + a_{\rm AMOC}^{i} \cdot (\Psi_{\rm AMOC}^{\prime})^{\rm model}(t), \tag{4}$$

and minimizing $J = \sum_{t_{oc}=t_{st}}^{t_{end}} [(T'_{oc})^{fit}(t_{oc}) - (T'_{oc})^{model}(t_{oc})]^2$. This results in estimates of the coefficients a_{Tatm}^i and a_{AMOC}^i (see Fig. S3 for examples). In this study, these coefficients are determined by an ensemble of 80 simulations performed with the B3D model, and by a deglaciation simulation using the LVC model²³ (Supplementary information).

 $T'_{\rm atm}$ is ideally global mean air temperature, because AMOC changes hardly affect global mean tempera-162 tures. Global mean $T'_{\rm atm}$ captures temperature changes induced by changes of greenhouse gases and surface 163 albedo. Unfortunately, to this date there is no direct paleoclimate proxy of global mean air temperature. 164 It is therefore determined from an Antarctic air temperature reconstruction as described in the main text. 165 Alternatively to global mean air temperatures, Antarctic air temperatures could be used, because $T'_{\rm atm}$ 166 does not necessarily need to be independent of variations in AMOC (Fig. S7). It is important that in this 167 case also Antarctic air temperatures are used for the determination of the coefficients of the reconstruction 168 scheme. However, a prerequisite for this is that the polar amplification must be realistically represented in 169 the model where the coefficients are determined. Furthermore, Antarctic air temperatures may addition-170 ally be influenced by quantities other than the AMOC which can lead to a systematic bias of the AMOC 171 reconstruction. 172

¹⁷³ Solving Eq. (1) for Ψ'_{AMOC} is trivial if only one SST time series reconstruction is used. If more than one $T_{oc}^{\prime n}$ ¹⁷⁴ reconstruction is available, a least squares solution can be calculated:

$$\Psi'_{\rm AMOC}(t) = \sum_{i=1}^{n} \frac{w_i}{a^i_{\rm AMOC}} [T'^i_{\rm oc}(t) - a^i_{\rm Tatm} T'_{\rm atm}(t)],$$

where w_i is the weight of location *i* and is determined by the least squares solver. Note that inaccurate coefficients a_{Tatm}^i may lead to a trend in Ψ'_{AMOC} in the case where T'_{oc}^i and T'_{atm} change simultaneously such as during a deglaciation. This is illustrated in a hindcast of the LVC deglaciation simulation where all a_{Tatm}^i are set to 0 K K⁻¹ (Fig. S8).

Reconstructions of sea-surface and air temperatures are generally not equally spaced in time. Hence, the time series are splined using a cutoff period of τ_{cut} years²⁹. τ_{cut} must be chosen appropriately depending on the resolution of the selected atmosphere and SST time series. A smaller τ_{cut} will result in a better resolved Ψ'_{AMOC} reconstruction.

The sensitivity of SSTs to changes in Ψ'_{AMOC} depends on the location. Some locations are not, or only weakly sensitive to changes in Ψ'_{AMOC} . At these locations, a_{AMOC} is close to zero. This may lead to large uncertainties in the AMOC reconstruction, because in this case Eq. (2) approaches singularity. However, as long as the suite of T'_{oc}^{i} time series includes data at locations that are sensitive to circulation changes, the least squares solution of the system of equations will assign very small weights to the time series that are only weakly sensitive to AMOC.

189 Data

- 190 SST and air temperature data used in the reconstruction (Table S1) were obtained at http://www.pangaea.
- de, http://www.ncdc.noaa.gov/paleo, or by personal communication. See Table S2 for specific links.

¹⁹² Additional information

¹⁹³ Supplementary data and information are available in the online version of the paper. Correspondence and ¹⁹⁴ requests for materials should be addressed to S.P.R.

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Author contributions

S.P.R. and T.F.S. developed the method, A.T. contributed with ideas on method verification and AMOC fingerprinting, S.P.R. performed the B3D simulations, L.M. provided the LVC simulation, J.O.G. provided support of the sediment core data, and S.P.R. and T.F.S. wrote the manuscript. All authors discussed the results and provided input on the manuscript.

²⁰⁴ Competing financial interests

²⁰⁵ The authors declare no competing financial interests.

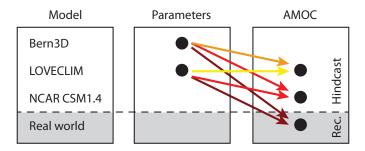


FIG. 1. Overview of the performed AMOC hindcasts and reconstructions. The color of the arrows indicates the order of appearance within the paper (going from light to dark shading). The center column indicates which models are used to calculate the coefficients of the reconstruction scheme. For example, coefficients determined from the Bern3D model are used to hindcast AMOC in LOVECLIM and NCAR models for verification with the directly simulated AMOC in these models, and to reconstruct AMOC from paleoclimate records.

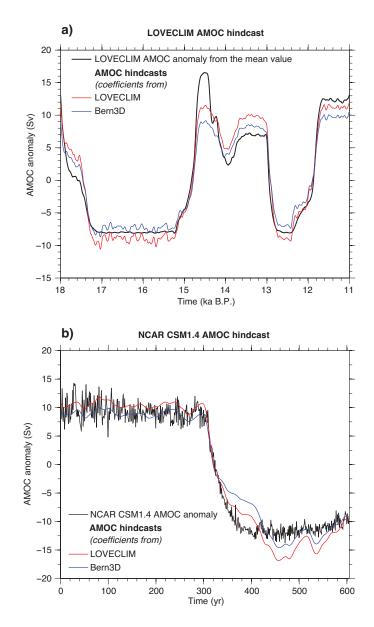


FIG. 2. AMOC hindcasts of two model simulations. The hindcast is given as anomaly from the time series mean value. a) Hindcast of the AMOC of a LOVECLIM deglaciation simulation²³ (Fig. S4) using 14 Atlantic SST time series (Fig. S1). Black line: LOVECLIM AMOC anomaly. Red line: AMOC hindcast calculated from coefficients determined by the LOVECLIM simulation. Blue line: Hindcast from coefficients determined by Bern3D. The coefficients are given in Tables S3 and S4. b) Hindcast of the AMOC of a freshwater forcing simulation²⁴ (Fig. S5). Black line: NCAR CSM1.4 AMOC anomaly. Red line: AMOC hindcast calculated from coefficients determined by LOVECLIM. Blue line: Hindcast from coefficients determined by Bern3D.

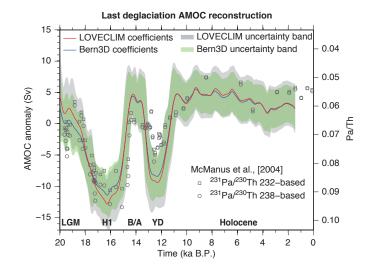


FIG. 3. Reconstruction of the AMOC of the last deglaciation. The reconstruction is given as anomalies from the time series mean value. For the reconstruction, 13 Atlantic SST time series are used (Fig. S1; Table S1) and global mean air temperature, which is determined from the Antarctic air temperature reconstruction from the EPICA Dome C Antarctic ice core²⁵ by applying a polar amplification factor and lowpass filtering the data to remove local-to-hemispheric scale variability. Blue line and green uncertainty band: Reconstruction where the coefficients calculated from Bern3D are used. Red line and gray uncertainty band: Reconstruction where the coefficients calculated from the LOVECLIM model are used. The uncertainty bands are determined by uncertainty propagation of the age-scale uncertainties of ± 500 years, and the uncertainties that arise from the ocean and air temperature time series reconstructions. The uncertainty band of the reconstruction based on Bern3D coefficients additionally includes coefficient uncertainties. The reconstruction shows different AMOC regimes during the Last Glacial Maximum (LGM), Heinrich stadial 1 (H1), the Bølling/Allerød (B/A), the Younger Dryas (YD) stadial, and during the Holocene. Dark gray symbols: Qualitative AMOC reconstruction from sedimentary ²³¹Pa/²³⁰Th².

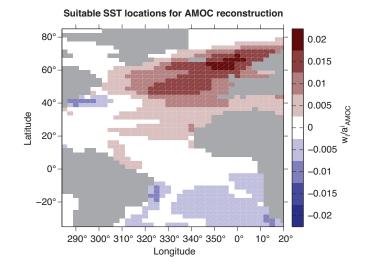


FIG. 4. Suitable locations of SST records for the AMOC reconstruction. An AMOC hindcast of the NCAR CSM1.4 simulation is calculated based on SST time series of every LOVECLIM surface Atlantic model cell and by using reconstruction scheme coefficients determined from the LOVECLIM model. w_i/a_{AMOC}^i is the spatial fingerprint of the AMOC (Eq. 3). It determines the influence of every location *i* on the AMOC hindcast. The larger the absolute values (the darker the colors), the larger the influence of the location on the AMOC hindcast.

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